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FCC Physics Prospects

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Why don't we see as yet the new physics we expected to be present around the TeV scale ?

- Is the mass scale beyond the LHC reach ?
- Is the mass scale within LHC's reach, but final states are elusive to the direct search ?

These two scenarios are a priori equally likely, but they impact in different ways the future of HEP, and thus the assessment of the physics potential of possible future facilities

Readiness to address both scenarios is the best hedge for the field:

- precision \Rightarrow higher statistics, better detectors and experimental conditions
- sensitivity (to elusive signatures) \Rightarrow ditto
- •extended energy/mass reach ⇒ higher energy

The physics potential (the "case") of a future facility for HEP should be weighed against criteria such as:

(1) the guaranteed deliverables:

 knowledge that will be acquired independently of possible discoveries (the value of "measurements")

(2) the **exploration potential:**

- target broad and well justified BSM scenarios but guarantee sensitivity to more exotic options
- exploit both direct (large Q^2) and indirect (precision) probes
- (3) the potential to provide conclusive **yes/no answers** to relevant, broad questions.

Future Circular Collider



What a future circular collider can offer

- <u>Guaranteed deliverables</u>:
 - study of Higgs and top quark properties, and exploration of EWSB phenomena, with the best possible precision and sensitivity
- Exploration potential:
 - \bullet exploit both direct (large Q²) and indirect (precision) probes
 - enhanced mass reach for direct exploration at 100 TeV
 - E.g. match the mass scales for new physics that could be exposed via indirect precision measurements in the EW and Higgs sector
- <u>Provide firm Yes/No answers</u> to questions like:
 - is there a TeV-scale solution to the hierarchy problem?
 - is DM a thermal WIMP?
 - could the cosmological EW phase transition have been 1st order?
 - could baryogenesis have taken place during the EW phase transition?
 - could neutrino masses have their origin at the TeV scale?

Event rates: examples

н	Ζ	W	t	т(←Z)	b(←Z)	c(←Z)
10 ⁶	5 10 ¹²	10 ⁸	10 ⁶	3 10 ¹¹	1.5 10 ¹²	2 10 ¹²
	н	b	t	W(·	←t)	r(←W←t)
2.5	10 ¹⁰	10 ¹⁷	10 ¹²	10	12	10 ¹¹
h		н			t	
		2.5 10 ⁶			2 10 ⁷	
	H 10 ⁶ 2.5	Η Ζ 10 ⁶ 5 10 ¹² Η 2.5 10 ¹⁰	H Z W 106 5 1012 108 H b 1017 2.5 1010 H H H H H 1017 1017 H H H H H H 1017 1017 H H H <t< td=""><td>H Z W t 106 5 1012 108 106 H b t 2.5 1010 1017 1012 H H H H 1017 1012 1012 H H H H H H H H 1017 1012 1012 H H H H H H H H H H H H H H H H H H H H</td><td>H Z W t $\tau(\leftarrow Z)$ 106 5 1012 108 106 3 1011 H b t W(4) 2.5 1010 1017 1012 10 H</td><td>H Z W t $\tau(\leftarrow Z)$ $b(\leftarrow Z)$ 106 5 1012 108 106 3 1011 1.5 1012 H b t W(\leftarrowt) H W(\leftarrowt) H</td></t<>	H Z W t 106 5 1012 108 106 H b t 2.5 1010 1017 1012 H H H H 1017 1012 1012 H H H H H H H H 1017 1012 1012 H H H H H H H H H H H H H H H H H H H H	H Z W t $\tau(\leftarrow Z)$ 106 5 1012 108 106 3 1011 H b t W(4) 2.5 1010 1017 1012 10 H	H Z W t $\tau(\leftarrow Z)$ $b(\leftarrow Z)$ 106 5 1012 108 106 3 1011 1.5 1012 H b t W(\leftarrow t) H W(\leftarrow t) H

(1) guaranteed deliverables: Higgs properties

https://arxiv.org/pdf/1708.08912.pdf

	Model	$b\overline{b}$	$c\overline{c}$	gg	WW	au au	ZZ	$\gamma\gamma$	$\mu\mu$
1	MSSM [40]	+4.8	-0.8	- 0.8	-0.2	+0.4	-0.5	+0.1	+0.3
2	Type II 2HD $[42]$	+10.1	-0.2	-0.2	0.0	+9.8	0.0	+0.1	+9.8
3	Type X 2HD $[42]$	-0.2	-0.2	-0.2	0.0	+7.8	0.0	0.0	+7.8
4	Type Y 2HD $[42]$	+10.1	-0.2	-0.2	0.0	-0.2	0.0	0.1	-0.2
5	Composite Higgs [44]	-6.4	-6.4	-6.4	-2.1	-6.4	-2.1	-2.1	-6.4
6	Little Higgs w. T-parity $[45]$	0.0	0.0	-6.1	-2.5	0.0	-2.5	-1.5	0.0
7	Little Higgs w. T-parity $[46]$	-7.8	-4.6	-3.5	-1.5	-7.8	-1.5	-1.0	-7.8
8	Higgs-Radion $[47]$	-1.5	- 1.5	+10.	-1.5	-1.5	-1.5	-1.0	-1.5
9	Higgs Singlet $[48]$	-3.5	-3.5	-3.5	-3.5	-3.5	-3.5	-3.5	-3.5

5 – 10 %

> 10%

NB: when the b coupling is modified, BR deviations are smaller than the square of the coupling deviation. Eg in model 5, the BR to b, c, tau, mu are practically SM-like

(sub)-% precision must be the goal to ensure 3-5σ evidence of deviations, and to cross-correlate coupling deviations across different channels

The absolutely unique power of $e^+e^- \rightarrow ZH$ (circular or linear):

- the model independent absolute measurement of HZZ coupling, which allows the subsequent:
 - sub-% measurement of couplings to W, Z, b, T
 - % measurement of couplings to gluon and charm



 $p(H) = p(e^-e^+) - p(Z)$

=> [p(e⁻e⁺) – p(Z)]² peaks at m²(H)

reconstruct Higgs events independently of the Higgs decay mode!



<u>The absolutely unique power of pp \rightarrow H+X:</u>

- the extraordinary statistics that, complemented by the per-mille e^+e^- measurement of eg BR(H \rightarrow ZZ*), allows
 - the sub-% measurement of rarer decay modes
 - the ~5% measurement of the Higgs trilinear selfcoupling
- the huge dynamic range (eg pt(H) up to several TeV), which allows to
 probe d>4 EFT operators up to scales of several TeV
 - search for multi-TeV resonances decaying to H, or extensions of the Higgs sector

	gg→H	VBF	WH	ZH	ttH	нн
N ₁₀₀	24 x 10 ⁹	2.1 x 10 ⁹	4.6 x 10 ⁸	3.3 x 10 ⁸	9.6 x 10 ⁸	3.6 x 10 ⁷
N100/N14	180	170	100	110	530	390

 $N_{100} = \sigma_{100 \text{ TeV}} \times 30 \text{ ab}^{-1}$

 $N_{14} = \sigma_{14 \text{ TeV}} \times 3 \text{ ab}^{-1}$

H at large рт



- Hierarchy of production channels changes at large p_T(H):
 - $\sigma(ttH) > \sigma(gg \rightarrow H)$ above 800 GeV
 - $\sigma(VBF) > \sigma(gg \rightarrow H)$ above 1800 GeV

Three kinematic regimes

- Inclusive production, $p_T > 0$:
 - largest overall rates
 - most challenging experimentally:
 - triggers, backgrounds, pile-up \Rightarrow low efficiency, large systematics
 - \blacksquare det simulations challenging, likely unreliable \Rightarrow regime not studied so far

• <u>p⊤ ≳ 100 GeV :</u>

- stat uncertainty ~few × 10⁻³ for $H \rightarrow 4I, \gamma\gamma, ...$
- improved S/B, realistic trigger thresholds, reduced pile-up effects ?
- current det sim and HL-LHC extrapolations more robust
- ➡ focus of FCC CDR Higgs studies so far
- sweet-spot for precision measurements at the sub-% level

• <u>p⊤ ≳ TeV :</u>

- stat uncertainty O(10%) up to 1.5 TeV (3 TeV) for $H \rightarrow 4I$, $\gamma\gamma$ ($H \rightarrow bb$)
- new opportunities for reduction of syst uncertainties (TH and EXP)
- different hierarchy of production processes
- indirect sensitivity to BSM effects at large Q² , complementary to that emerging from precision studies (eg decay BRs) at Q~m_H

$gg \rightarrow H \rightarrow \gamma \gamma$ at large p_T



lacksquare	At LHC, S/B	in the $H \rightarrow \gamma \gamma$ channel is O(few %)	

- At FCC, for p_T(H)>300 GeV, S/B~I
- Potentially accurate probe of the H pt spectrum up to large pt

δ _{stat}	р _{т,min} (GeV)
0.2%	100
0.5%	400
1%	600
10%	1600



Normalize to BR(4I) from ee => sub-% precision for absolute couplings

Future work: explore in more depth data-based techniques, to <u>validate and</u> <u>then reduce</u> the systematics in these ratio measurements, possibly moving to lower pt's and higher stat



Higgs couplings after FCC-ee / hh

	HL-LHC	FCC-ee	FCC-hh
δΓΗ / ΓΗ (%)	SM	1.3	tbd
δg _{HZZ} / g _{HZZ} (%)	1.5	0.17	tbd
δднww / днww (%)	1.7	0.43	tbd
δg _{Hbb} / g _{Hbb} (%)	3.7	0.61	tbd
δg_{Hcc} / g_{Hcc} (%)	~70	1.21	tbd
δg _{Hgg} / g _{Hgg} (%)	2.5 (gg->H)	1.01	tbd
δgнττ / gнττ (%)	1.9	0.74	tbd
δg _{Hµµ} / g _{Hµµ} (%)	4.3	9.0	0.65 (*)
δg _{Hγγ} / g _{Hγγ} (%)	1.8	3.9	0.4 (*)
δg _{Htt} / g _{Htt} (%)	3.4	~10 (indirect)	0.95 (**)
δg _{HZγ} / g _{HZγ} (%)	9.8	_	0.9 (*)
δдннн / дннн (%)	50	~44 (indirect)	3.5
BR _{exo} (95%CL)	BR _{inv} < 2.5%	< 1%	BR _{inv} < 0.025%

NB

BR(H→ZY,YY) ~O(10⁻³) ⇒ O(10⁷) evts for Δ_{stat} ~% BR(H→µµ) ~O(10⁻⁴) ⇒ O(10⁸) evts for Δ_{stat} ~%



pp collider is essential to beat the % target, since no proposed ee collider can produce more than O(10⁶) H's

* From BR ratios wrt $B(H \rightarrow ZZ^*)$ @ FCC-ee

** From pp \rightarrow ttH / pp \rightarrow ttZ, using B(H \rightarrow bb) and ttZ EW coupling @ FCC-ee

The Higgs self-coupling at FCC-hh https://arxiv.org/abs/2004.03505



-2∆ In L

Figure 13. Expected negative log-Likelihood scan as a function of the trilinear self-coupling modifier $\kappa_{\lambda} = \lambda_3 / \lambda_3^{\text{SM}}$ in all channels, and their combination. The solid line corresponds to the scenario II for systematic uncertainties. The band boundaries represent respectively scenario I and III. The dashed line represents the sensitivity obtained including statistical uncertainties only, under the assumptions of scenario I.

Syst scenarios

		-		
	@68% CL	scenario I	scenario II	scenario III
s	stat only	2.2	2.8	3.7
δ_{μ}	stat + syst	2.4	3.5	5.1
$\delta_{\kappa_{\lambda}}$	stat only	3.0	4.1	5.6
	stat + syst	3.4	5.1	7.8

Table 7. Combined expected precision at 68% CL on the di-Higgs production cross- and Higgs self coupling using all channels at the FCC-hh with $\mathcal{L}_{int} = 30 \text{ ab}^{-1}$. The symmetrized value $\delta = (\delta^+ + \delta^-)/2$ is given in %.

- I. Target det performance: LHC Run 2 conditions
- II. Intermediate performance
- III. Conservative: extrapolated HL-LHC performance, with today's algo's (eg no timing, etc)

Expected precision on the Higgs self-coupling as a function of the integrated luminosity.

3-5 ab⁻¹ are sufficient to get below the 10% level

=> within the reach of the first 5yrs of FCC-hh running, in

the <u>"low" luminosity / low pileup</u> phase

=> the 10% precision threshold can be reached within the timescale of a similar measurement by CLIC @ 3 TeV

Extracting Higgs self-coupling from HH at FCC:

the power of ee/hh synergy & complementarity

At FCC-hh we can precisely measure HH rate ... but, to interpret this as H selfcoupling:



Direct measurement of ttH coupling: from $R_t = \sigma(ttH)/\sigma(ttZ)$

FCC-hh can measure R_t with $\Delta R_t/R_t < 2\%$ but:



Unique at FCC-ee: $e^+e^- \rightarrow H$

D.D'Enterria et al, arXiv:2107.02686



Table 6. Individual significances (in std. deviations σ) expected per decay channel for *s*-channel Higgs boson production in e^+e^- collisions at FCC-ee for $\mathcal{L}_{int} = 10 \text{ ab}^{-1}$ and $\delta_{\sqrt{s}} = 4.1 \text{ MeV}$. The last column quotes the combined significance.

${\rm H} \to gg$	$\mathrm{H} \to \mathrm{WW}^* \to \ell \nu \ 2j; \ 2\ell \ 2\nu; \ 4j$	$\mathrm{H} \to \mathrm{ZZ}^* \to 2j \; 2\nu; \; 2\ell \; 2j; \; 2\ell \; 2\nu$	${\rm H} \to b \overline{b}$	$\mathrm{H} \to \tau_{\mathrm{had}} \tau_{\mathrm{had}}; c \overline{c}; \gamma \gamma$	Combined
1.1σ	$(0.53\otimes 0.34\otimes 0.13)\sigma$	$(0.32\otimes 0.18\otimes 0.05)\sigma$	0.13σ	$< 0.02\sigma$	1.3σ

(1) guaranteed deliverables: EW&flavour observables

The absolutely unique power of **Circular** e⁺e⁻:

e+e- → Z	e+e- → WW	т(←Z)	b(←Z)	c(←Z)	e+e- → tt
5 10 ¹²	10 ⁸	3 10 ¹¹	1.5 10 ¹²	10 ¹²	10 ⁶

=> O(10⁵) larger statistics than LEP at the Z peak and WW threshold

Flavour statistics from Z decays:

Particle production (10^9)	B^0	B^-	B_s^0	Λ_b	$c\overline{c}$	$\tau^- \tau^+$
Belle II	27.5	27.5	n/a	n/a	65	45
FCC-ee	400	400	100	100	800	220

Additional bonus wrt B factory: (i) Lorentz boost (ii) B hadrons not accessible at the Y(4S,5S) thresholds

EW parame **@ FCC-е**

Improvement wrt current to uncertainties:

- stat precision ~ 10-1000
- with exptl syst ~ > 10-50

Currently limited by TH sys =>

ee goals set during the on Workshop

	Observable	present value ± error	FCC-ee stat.	FCC-ee syst.
N parameters	m _Z (keV)	91186700±2200	5	100
@ FCC-ee	$\Gamma_{\rm Z}$ (keV)	2495200±2300	8	100
	R_l^Z (×10 ³)	20767±25	0.06	0.2-1.0
	α_{s} (m _Z) (×10 ⁴)	1196±30	0.1	0.4-1.6
	R_{b} (×10 ⁶)	216290±660	0.3	<60
	$\sigma_{\rm had}^0$ (×10 ³) (nb)	41541±37	0.1	4
	N_{ν} (×10 ³)	2991±7	0.005	1
rovement wrt current total ertainties:	$\sin^2 \theta_W^{eff}$ (×10 ⁶)	231480±160	3	2-5
at precision $\sim 10-1000$ smaller	$1/\alpha_{QED}(m_Z)$ (×10 ³)	128952±14	4	Small
ith exptl syst $\sim > 10-50$ smaller	$A_{\rm FB}^{b,0}$ (×10 ⁴)	992±16	0.02	1-3
rently limited by TH systematics	$A_{\rm FB}^{{\rm pol}, \tau}$ (×10 ⁴)	1498±49	0.15	<2
poals set during the ongoing	m _W (MeV)	80350±15	0.6	0.3
rkshop	$\Gamma_{\rm W}$ (MeV)	2085±42	1.5	0.3
	$\alpha_s (m_W) (\times 10^4)$	1170±420	3	Small
	$N_{\nu}(\times 10^3)$	2920±50	0.8	Small
	m _{top} (MeV)	172740±500	20	Small
	$\Gamma_{\rm top}$ (MeV)	1410±190	40	Small
	$\lambda_{\rm top}/\lambda_{\rm top}^{\rm SM}$	1.2±0.3	0.08	Small
crucial for ttH and HHH couplings at FCC-hh	ttZ couplings	±30%	▶ 0.5 - 1.5%	Small

Flavour probes: eg lepton universality in tau decays



Lorentz boost crucial!

			t metime [i	2]		
	Observable	Measurement	Current precision	FCC-ee stat.	Possible syst.	Challenge
	m _τ [MeV]	Threshold / inv. mass endpoint	1776.86 ± 0.12	0.004	0,04-0,1	Mass scale
	→ τ _τ [fs]	Flight distance	290.3 ± 0.5 fs	0.001	0.04	Vertex detector alignment
	Β(τ→evv) [%]	Selection of τ⁺τ,	17.82 ± 0.05	0.0001	0.000	Efficiency, bkg,
	Β(τ→μνν) [%]	state	17.39 ± 0.05	0.0001	0.003	Particle ID

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mtop: the advantage of circular

Optimisation of the threshold scan



23

Precision W physics at FCC-hh: LHC docet

ATLAS 2020: <u>arXiv:2007.14040</u>

CMS 2022: arXiv:2201.07861



^{~ 300} x HL-LHC statistics

(2) Direct discovery reach at high mass: the power of 100 TeV

ATLAS Preliminary

ATLAS SUSY Searches* - 95% CL Lower Limits

Moren 2014	$v_{\rm m} = 12 \text{ TeV}$
Model Signature ∫£ dr [16 ⁻¹] Mass limit Ref	erence
$\hat{q}\hat{q}, \hat{q} \rightarrow \hat{q}\hat{\ell}_{1}^{0}$ 0 c, μ 2-6 jets E_{T}^{vin} 36.1 (2x, 8x Degen.) 0.9 1.55 m($\hat{\ell}_{1}^{0}$)<100 GeV 11 mono-jet 1-3 jets E_{T}^{vin} 36.1 (1x, 8x Degen.) 0.43 0.71 m($\hat{\mu}_{1}$)<100 GeV 11 m($\hat{\mu}_{1}$)<100 GeV 11	12.02332
$g_{3}, g \to q \bar{q} \tilde{T}_{1}^{0}$ 0 c, μ 2-6 jets E_{7}^{min} 36.1 k 2.0 mt \tilde{C}_{1} 2.0 mt \tilde{C}_{1} 2.00 GeV 17	12.02332
$\frac{c}{k_{1}^{2}, k_{2}^{2} \rightarrow k_{1}^{2}(\ell) \tilde{k}_{1}^{0}} \qquad 3 e, \mu 4 \text{ jots} 3e, 1 k \frac{1.85}{2} \qquad \frac{1.85}{1.85} \qquad m(\tilde{k}_{1}^{2} \Rightarrow 00 \text{ GeV}) 1$	06.03731
$e_{r,\mu\mu} = 2 \text{ jets} E_{r}^{rav} = 36.1 \hat{x} = 1.2 \qquad \text{m}(\hat{r}) = 50 \text{ GeV} = 19$	05.11381
$3r_{\mu} 4 \text{ jets} 3s.1 \mathbf{x} 0.98 \mathbf{m}(\mathbf{x}) = \mathbf{m}(\mathbf{x}$	06.03731
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	CONF-2018-041 06.03731
$ \begin{array}{c c c c c c c c c c c c c c c c c c c $	266, 1711.03301 108.09266 106.03731
$ \frac{5}{6} \delta_1 \delta_1, \delta_1 \rightarrow \delta_2^{\mu} \rightarrow b \delta_1^{\mu} $ $ 0 c_1 \mu 6 \ b \ h_T^{\min} 139 \frac{5}{6} Forbidden 0.23 - 0.35 dm(\ell_1^0, \ell_1^0) - 100 \ GeV, m(\ell_1^0, \ell_1^0) - 100 \ GeV, m(\ell_1^0, \ell_1^0) - 100 \ GeV, m(\ell_1^0, \ell_1^0) - 100 \ GeV SU $	5Y-2018-31 5Y-2018-31
$\frac{1}{6} \tilde{I}_1 \tilde{I}_1, \tilde{I}_1 \rightarrow W h \tilde{k}_1^0 \text{ or } \tilde{k}_1^0 = 0.2 \text{ (e.m. 0.2 jets/1.2 b } L_T^{min} = 36.1 \tilde{I}_1 = 0.2 \text{ (e.m. 0.2 jets/1.2 b } L_T^{min} = 36.1 \tilde{I}_1 = 0.2 \text{ (f. 0.1 min)} = 0$	09.04183, 1711.11520
$\frac{\delta_{1}}{\tau_{1}} \tilde{f}_{1} \tilde{f}_{1}, \text{ Well-Tempered LSP} \qquad Multiple \qquad 36.1 \tilde{f}_{1} \qquad \qquad 0.48-0.84 \qquad m(\tilde{t}_{1}^{2})-150 \text{ GeV}, n(\tilde{t}_{1}^{2})-5 \text{ GeV}, \tilde{t}_{1} \times \tilde{t}_{2} \qquad 1709.047$	183, 1711.11520
$\frac{1}{6} f_1 f_1, f_1 \rightarrow \hat{\tau}_1 bv, \hat{\tau}_1 \rightarrow \hat{\tau}_G $ $\frac{1}{7} + 1 e_{\mu,\tau} 2 jots' 1 b E_{\tau}^{rres} 38.1 I_1 $ $\frac{1}{16} $ $\frac{1}{16} $ $\frac{1}{16} $ $\frac{1}{16} $	03.10178
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	05.01649
$0 e_{\mu}$ mono-jet E_T^{mino} 36.1 \bar{I}_1 0.43 $m(\bar{I}_1)=5 GeV$ 17	11.03301
$\tilde{t}_{2}\tilde{t}_{2}, \tilde{t}_{2} \rightarrow \tilde{t}_{1} + h \qquad 1 \cdot 2 \epsilon \mu \qquad 4 \ b \qquad E_{T}^{\pi i m} 36.1 \qquad \tilde{t}_{2} \qquad 0.32 \cdot 0.00 \qquad m(\tilde{t}_{1}^{2}) - 0 \ 6 \epsilon V, \\ m(\tilde{t}_{1}) - m(\tilde{t}_{1}^{2}) - 180 \ G \epsilon V \qquad 12 \cdot 100 \ G \epsilon V \ 12 \cdot 100 \ G$	36603986
$ \hat{\chi}_{1}^{*} \hat{\chi}_{2}^{0} \text{ via } WZ \qquad \begin{array}{cccccccccccccccccccccccccccccccccc$	94, 1806.02293 12.08119
$k_1^{\pm} k_1^{\pm} \text{ via } WW = 2 c_{\mu} \mu - k_T^{min} = 139 - k_1^{\pm} = 0.42$ ATLAS4	CONF-2019-008
$\hat{x}_{1}^{-}\hat{x}_{2}^{0}$ via W/h 0-1 e, μ 2 h h_{T}^{\min} 36.1 $\hat{x}_{1}^{+}/\hat{x}_{2}^{0}$ 0.68 m(\hat{v}_{1})=0 1/2	12.09432
$\frac{1}{5}$ $\hat{\chi}_1 \hat{\chi}_1 = \frac{1}{10}$ $\hat{\chi}_1^{(1)} = \frac{1}{10}$	CONF-2019-005
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	08.07875 08.07875
	CONF-2019-008
$2 c_{,\mu} \ge 1 - k_T^{riv} = 36.1 - \frac{7}{2} - 0.10$ $m(\tilde{r}) - m(\tilde{r}_1^*) - 5 \text{ GeV} = 1$	12.08119
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	05.04030 04.03602
Direct $\hat{\chi}_1^+ \hat{\chi}_1^-$ prod., long-lived $\hat{\chi}_1^+$ Disapp. trk 1 jet E_T^{mix} 36.1 $\hat{\chi}_1^+$ 0.46 Pure Wino 11	12.02118 S.PUB.2017.019
Stable - R.hadron Multiple 98.1 - 20	51F 06 2017-010
Autor g relation Multiple 36.1 A 100	901,1808.04095
$V_{\rm ev} p_{\rm p} \rightarrow v_{\rm e} + x_{\rm e} v_{\rm p} \rightarrow p_{\rm e} \rightarrow p_{\rm e} + x_{\rm e} v_{\rm p} \rightarrow$	07.06079
$\frac{1}{2}, \frac{1}{2} \rightarrow agg$ 4-5 large- <i>R</i> jets 36.1 k [mit] $\frac{1}{2}$ [mit] \frac	04.03568
Multiple 36.1 2 [7] -2e-4, 2e-5] 1.05 2.0 m(1)-200 GeV, bino-800 ATLAS-4	CONF-2018-003
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	CONF-2018-003
$i_1i_1, i_1 \rightarrow b_3$ 2 jots + 2 b 36.7 i_1 (eq. b) 0.42 0.61	10.07171
$\frac{1}{1\mu} DV \qquad 136 \frac{1}{1\mu} \frac{0.4-1.45}{1.0} \qquad \frac{0.4-1.45}{1.0}$	10.05544 CONF-2019-006
nly a selection of the available mass limits on new states or 10 ⁻¹ Mass scale [TeV]	· · ·
molified models, c.f. refs, for the assumptions made.	
	0.4
	1.0
	7
	•

26

@100 TeV

Global EFT fits to EW and H observables at FCC-ee



Constraints on the coefficients of various EFT op's from a global fit of (i) EW observables, (ii) Higgs couplings and (iii) EW+Higgs combined. Darker shades of each color indicate the results neglecting all SM theory uncertainties.

s-channel resonances



100 TeV allow to directly access the mass scales revealed indirectly by precision EW and H measurements at the future e+e- factory

Matching this discovery reach with a lepton collider would require a multi-tens TeV facility (beyond-the-beyond?).

SUSY reach at 100 TeV



15-20 TeV squarks/gluinos would require a lepton collider in the ECM range of 30-50 TeV

(3) The potential for yes/no answers to important questions

WIMP DM theoretical constraints

For particles held in equilibrium by pair creation and annihilation processes, ($\chi \ \chi \leftrightarrow SM$)

$$\Omega_{\rm DM} h^2 \sim rac{10^9 {
m GeV}^{-1}}{M_{
m pl}} rac{1}{\langle \sigma v
angle}$$

For a particle annihilating through processes which do not involve any larger mass scales:

 $\langle \sigma v \rangle \sim g_{\rm eff}^4 / M_{\rm DM}^2$

$$\Omega_{\rm DM}h^2 \sim 0.12 \times \left(\frac{M_{\rm DM}}{2\,{\rm TeV}}\right)^2 \left(\frac{0.3}{g_{\rm eff}}\right)^4$$

$$\Omega_{wimp} h^2 \lesssim 0.12$$

$$M_{wimp} \lesssim 2\,{\rm TeV}\left(\frac{g}{0.3}\right)^2$$

K. Terashi, R. Sawada, M. Saito, and S. Asai, *Search for WIMPs with disappearing track signatures at the FCC-hh*, (Oct, 2018) . https://cds.cern.ch/record/2642474.

Disappearing charged track analyses (at ~full pileup)



=> coverage beyond the upper limit of the thermal WIMP mass range for both higgsinos and winos !!



The nature of the EW phase transition



Strong Ist order phase transition is required to induce and sustain the out of equilibrium generation of a baryon asymmetry during EW symmetry breaking

Strong Ist order phase transition $\Rightarrow \langle \Phi_C \rangle > T_C$

In the SM this requires $m_H \approx 80$ GeV, else transition is a smooth crossover.

Since $m_H = 125$ GeV, **new physics**, coupling to the Higgs and effective at **scales O(TeV)**, must modify the Higgs potential to make this possible

Probe higher-order terms of the Higgs potential (selfcouplings)
 Probe the existence of other particles coupled to the Higgs

Constraints on models with 1st order phase transition at the FCC

$$\begin{split} V(H,S) &= -\mu^2 \left(H^{\dagger} H \right) + \lambda \left(H^{\dagger} H \right)^2 + \frac{a_1}{2} \left(H^{\dagger} H \right) S \\ &+ \frac{a_2}{2} \left(H^{\dagger} H \right) S^2 + \frac{b_2}{2} S^2 + \frac{b_3}{3} S^3 + \frac{b_4}{4} S^4. \end{split}$$

Combined constraints from precision Higgs measurements at FCC-ee and FCC-hh **Direct detection of extra Higgs states at** FCC-hh



Remarks

- Apparently, adding the self-coupling constraint does not add much in terms of exclusion power, wrt the HZZ coupling measurement ...
- ... BUT, should HZZ deviate from the SM, λ_{HHH} is necessary to break the degeneracy among all parameter sets leading to the same HZZ prediction



- The concept of "which experiment sets a better constraint on a given parameter" is a very limited comparison criterion, which looses value as we move from "setting limits" to "diagnosing observed discrepancies"
- Likewise, it's often said that some observable sets better limits than others: "all known model predict deviations in X larger than deviations in Y, so we better focus on X". But once X is observed to deviate, knowing the value of Y could be absolutely crucial
- Redundancy and complementarity of observables is of paramount importance

MSSM Higgs @ 100 TeV



 N. Craig, J. Hajer, Y.-Y. Li, T. Liu, H. Zhang,
 J. Hajer, Y.-Y. Li, T. Liu, and J. F. H. Shiu,

 arXiv:1605.08744
 arXiv:1504.07617

... and much more ...

- Countless studies of discovery potential for multiple BSM scenarios, from SUSY to heavy neutrinos, from very low masses to very high masses, LLPs, DM, etcetcetc, with plenty of opportunities for direct discovery even at FCCee and FCC-eh
- Sensitivity studies to SM deviations in the properties of top quarks, flavour physics in Z decays: huge event rates offer unique opportunities, that cannot be matched elsewhere

• Operations with heavy ions: new domains open up at 100 TeV in the study of high-T/high-density QCD. Broaden the targets, the deliverables, extend the base of potential users, and increase the support beyond the energy frontier community

Final remarks

- The study of the SM will not be complete until we clarify the nature of the Higgs mechanism and exhaust the exploration of phenomena at the TeV scale: many aspects are still obscure, many questions are still open.
- The exptl program possible at a future circular collider facility, combining a versatile high-luminosity e⁺e⁻ circular collider, with a follow-up pp collider in the 100 TeV range, offers unmatchable breadth and diversity: concrete, compelling and indispensable Higgs & SM measurements enrich a unique direct & indirect discovery potential
- The next 5-6 years, before the next review of the European Strategy for Particle Physics, will be critical to reach the scientific consensus and political support required to move forward