

Symphony of Spacetime: Tuning into EMRI Resonances with LISA

Kyriakos Destounis, CENTRA-IST, Lisbon, Portugal

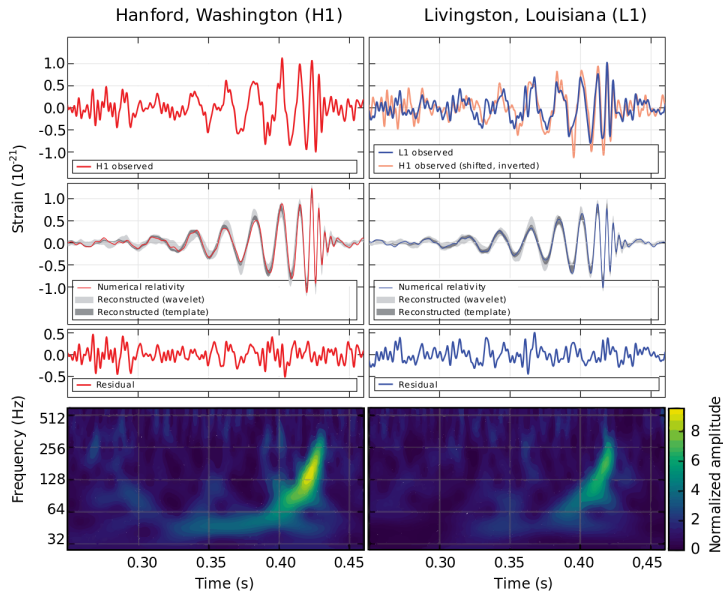
NEB-21, Corfu, Greece, 3 September, 2025



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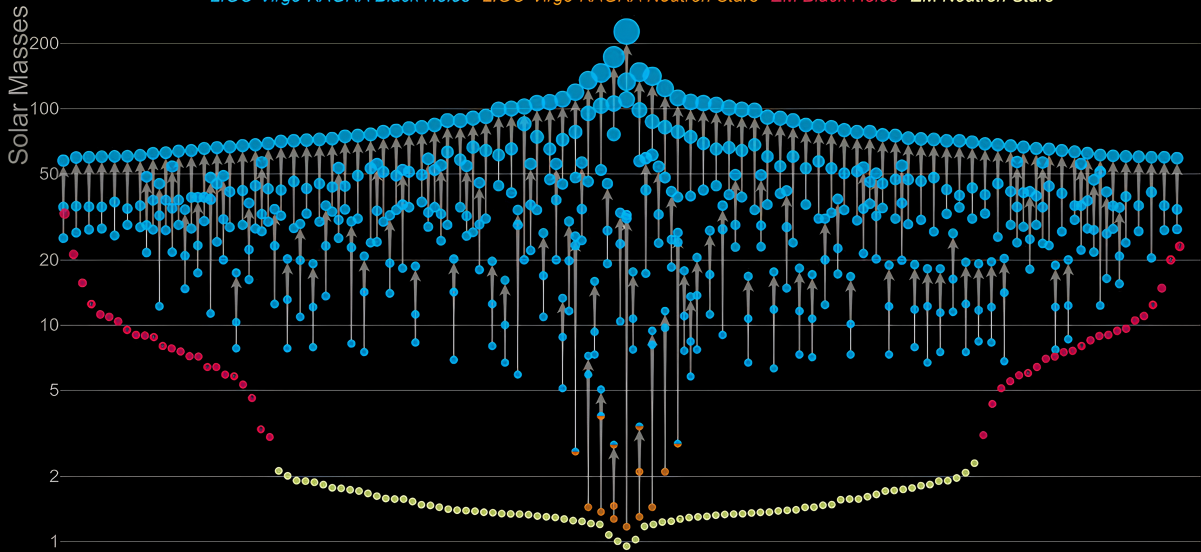
[LIGO and Virgo Collaboration, PRL (2016)]

LIGO/Virgo/KAGRA:

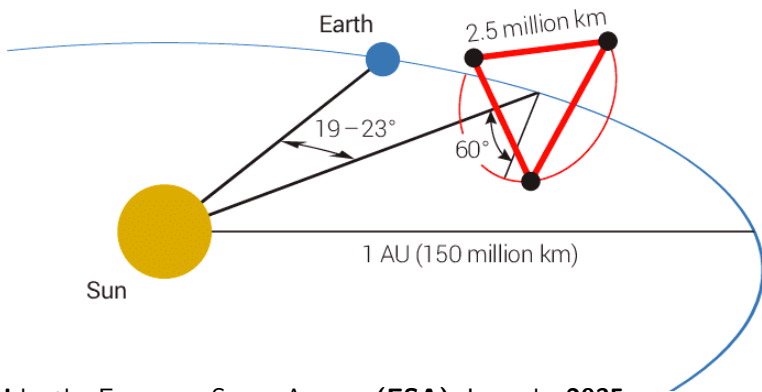
- ground-based detectors.
- 4th observation run.
- Increasing sensitivity.
- kHz band.
- 300+ detections.

Masses in the Stellar Graveyard

LIGO-Virgo-KAGRA Black Holes *LIGO-Virgo-KAGRA Neutron Stars* *EM Black Holes* *EM Neutron Stars*



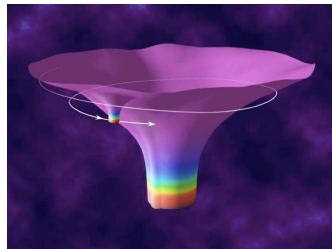
Laser Interferometer Space Antenna (LISA)



- **Adopted** by the European Space Agency (**ESA**). Launch: **2035**.
- **Multiple targets:** supermassive black-hole binaries, **extreme-mass-ratio inspirals**, compact binaries in the Milky Way, tests of GR, nature of black holes, environmental effects, dark matter, Λ CDM, stochastic gravitational-wave background, inflation.
- **Greece plays an active role (through the Δ ISA-GR Member Group) in the LISA Consortium.** [Karnesis+, IJMPD (2024)]

Extreme-mass-ratio inspirals

- **Extreme-mass-ratio inspirals (EMRIs):** primary supermassive black hole (M) + secondary stellar-mass compact object (μ).
- **Originate in** dense stellar clusters and **galactic cores**.
- The **mass ratio of EMRIs** (μ/M) span in the range $\sim 10^{-7} - 10^{-4}$.
- **Generate GWs** in the frequency range $\sim 10^{-4} - 10^{-1}$ Hz.
- Perform $\sim 10^5$ **revolutions in the strong-field regime**.
- **Visible for years** in the mHz band (**LISA**).
- **Rich phenomenology**.



EMRI modeling: adiabatic approximation + NK scheme

Short-timescale evolution: geodesics

Stationary/axisymmetric spacetime: $ds^2 = g_{tt}dt^2 + 2g_{t\phi}dtd\phi + g_{rr}dr^2 + g_{\theta\theta}d\theta^2 + g_{\phi\phi}d\phi^2$

Geodesic equations: $\ddot{x}^\kappa + \Gamma_{\lambda\nu}^\kappa \dot{x}^\lambda \dot{x}^\nu = 0$

Stationarity: $\mathcal{E} = -\mu (g_{tt}\dot{t} + g_{t\phi}\dot{\phi})$, **Axisymmetry:** $\mathcal{L}_z = \mu (g_{t\phi}\dot{t} + g_{\phi\phi}\dot{\phi})$

Carter constant $\leftrightarrow$ separation of radial and polar motion $\leftrightarrow$ **integrability**

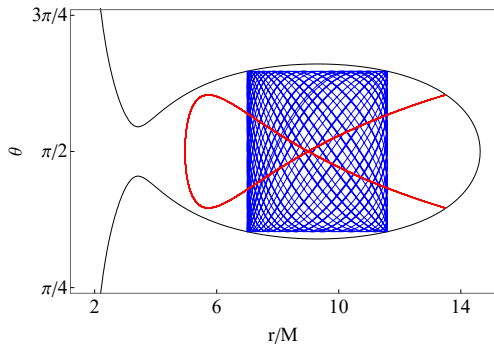
Non-integrability $\leftrightarrow$ **chaos** [Contopoulos, (2002)]

Long-timescale evolution: gravitational radiation reaction (made easy?)

Adiabatic approximation + numerical kludge (NK) scheme: Change in the momenta $(\mathcal{E}, \mathcal{L}_z, \mathcal{Q})$ is small over a single orbit $\rightarrow$ **approximate the dissipative fluxes** with post-Newtonian formulae [Glampedakis+, PRD (2002), Gair+, PRD (2008), Apostolatos+, PRL (2009)]

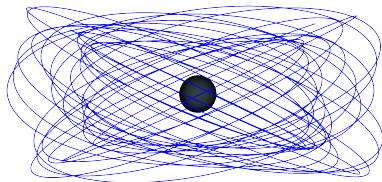
- **Short timescales:** orbit neglects radiative backreaction
- **Long timescales:** inspiral behaves like a ‘flow’ through successive geodesics

(r, θ) -bound motion

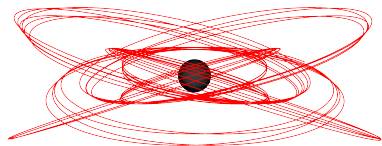


3D Cartesian coordinates

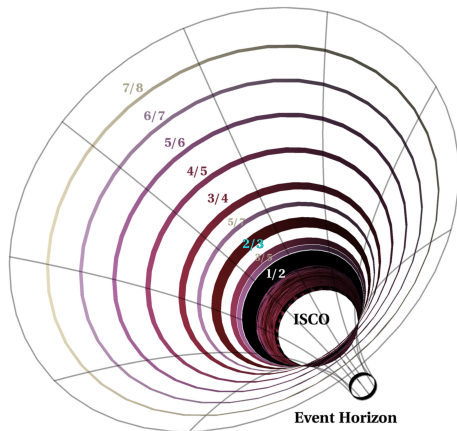
Generic orbit: $\omega_r/\omega_\theta = \text{irrational}$



Resonant orbit: $\omega_r/\omega_\theta = \text{rational}$



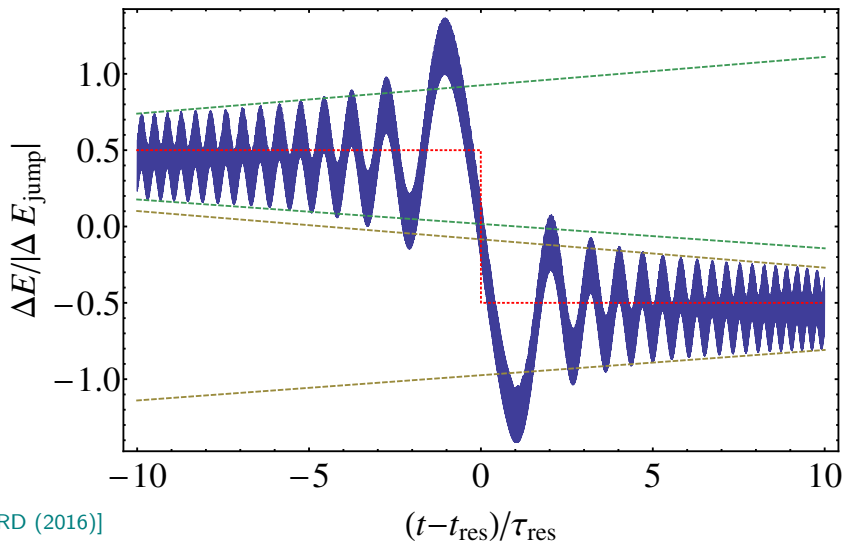
Resonances in Kerr EMRIs



[Brink, Geyer, Hinderer, PRL (2015)]

- Resonances are **ubiquitous** in EMRIs.
- Occur when **two orbital frequencies form a rational ratio**.
- Last for **dozens of cycles** in Kerr.
[Ruangsi+, PRD (2014), Brink+, PRD (2015)]
- Lead to **resonant “kicks”** [Flanagan, Hinderer, PRL (2012), Berry+, PRD (2016)]
- **Impact parameter estimation** [Berry+, PRD (2016), Speri+, PRD (2021)]

Resonant kick in fluxes of a Kerr EMRI (2 : 3)



[Berry+, PRD (2016)]

Effective resonance model: including resonant effects in NK

Since the **leading-order dissipative component** of the self-force alone is **insufficient to model the effect of resonances**, one must find a way to **enhance the NK fluxes**.

The **effective resonance model (ERM)** simulates the resonant contribution to the fluxes by adding to the NK the following **impulse function** $w(t)$: [\[Speri+, PRD \(2021\)\]](#)

$$\frac{dJ_i}{dt} = f_{\text{NK}} [1 + C_i w(t)], \quad w(t) = \begin{cases} \frac{1 + \cos \left[4\pi \left(\frac{t-t_0}{t_{\text{res}}} \right)^2 \right]}{\int_{-1/2}^{+1/2} 1 + \cos[4\pi x^2] dx}, & t \in [t_{\text{start}}, t_{\text{start}} + t_{\text{res}}], \\ 0, & \text{elsewhere,} \end{cases}$$

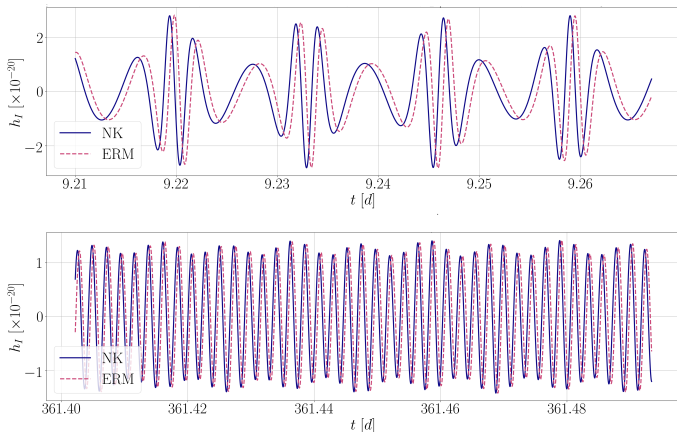
where $J_i = (\mathcal{E}, \mathcal{L}_z, \mathcal{Q})$, $C_i = (C_{\mathcal{E}}, C_{\mathcal{L}_z}, C_{\mathcal{Q}})$ are the **resonance coefficients**.

The **resonance duration** can be found by a Taylor expansion, resulting to

$$t_{\text{res}} = \sqrt{\frac{4\pi}{|l^* \dot{\omega}_{r_0} + m^* \dot{\omega}_{\theta_0}|}},$$

where $l^* \omega_{r_0} + m^* \omega_{\theta_0} = 0$ at $t = t_0$ is the **resonance condition**. [\[Ruangsri+, PRD \(2014\)\]](#)

Dephasing between NK and NK+ERM (2 : 3 resonance)

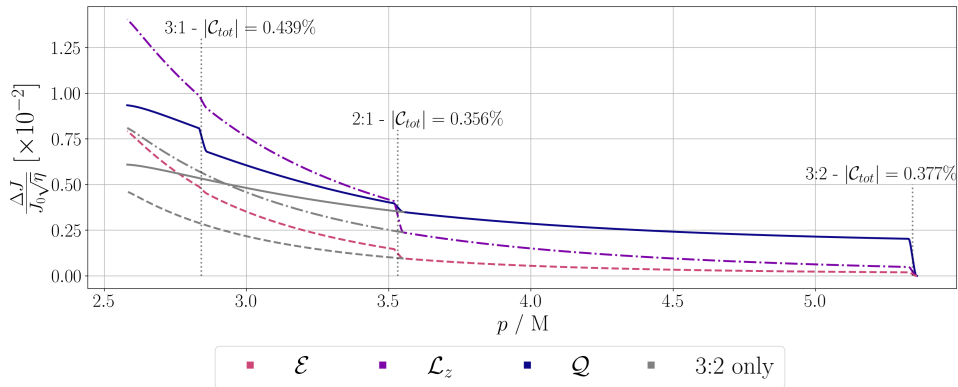


EMRI₁(top): $a = 0.80$, $p/M = 7.03$, $e = 0.40$, $\iota = 0.70$, $\mathcal{M}_{\text{top}} = 0.24$.

EMRI₂(bottom): $a = 0.80$, $p/M = 8.46$, $e = 0.20$, $\iota = 1.10$, $\mathcal{M}_{\text{bottom}} = 0.99$.

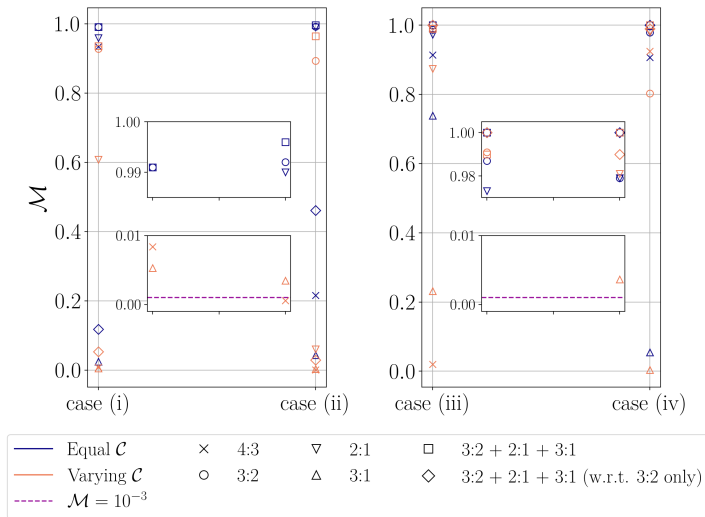
[Levati, Cardenas-Avendan˜o, KD, Pani, PRD (2025)]

Multiple resonance-crossing with ERM: fluxes

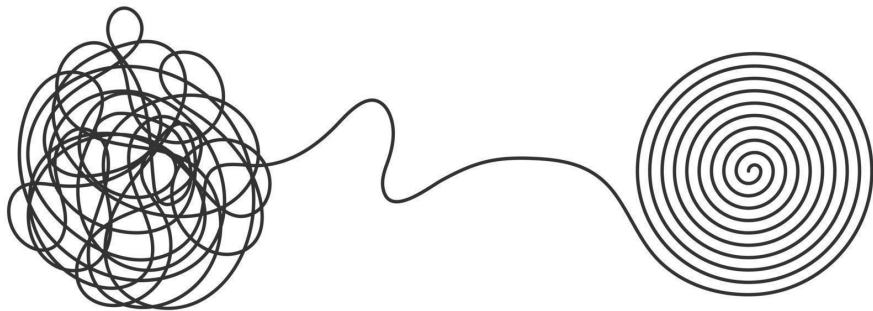


- Taking into account the (3 : 2) already affects the inspiral and resulting GW. [Flanagan+, PRL (2012), Speri+, PRD (2021)]
- Taking into account **sub-dominant resonances** affects the inspiral even more when compared with the inspiral that only includes (3 : 2). [Levati, Avendanõ, KD, Pani, PRD (2025)]

GW mismatch between NK and NK+multiple ERM



Chaos and Order



EMRIs and chaos: a non-Kerr primary

[KD, Suvorov, Kokkotas, PRD (2020)]

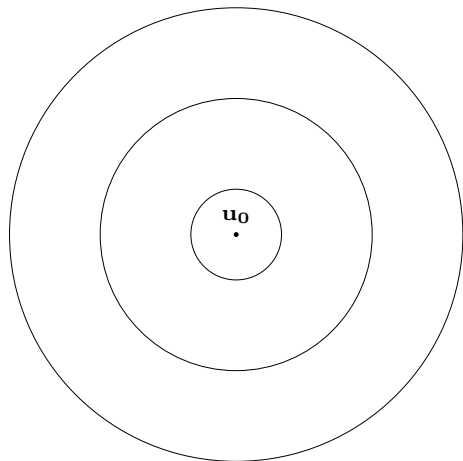
$$ds^2 = - \frac{\Sigma[(\alpha_Q/r) M^3 + \Delta - a^2 A(r)^2 \sin^2 \theta]}{[(r^2 + a^2) - a^2 A(r) \sin^2 \theta]^2} dt^2 - \frac{2a[(r^2 + a^2) A(r) - \Delta] \Sigma \sin^2 \theta}{[(r^2 + a^2) - a^2 A(r) \sin^2 \theta]^2} dt d\phi \\ + \frac{(\alpha_Q/r) M^3 + \Sigma}{\Delta} dr^2 + \Sigma d\theta^2 + \frac{\Sigma \sin^2 \theta [(r^2 + a^2)^2 - a^2 \Delta \sin^2 \theta]}{[(r^2 + a^2) - a^2 A(r) \sin^2 \theta]^2} d\phi^2,$$

$$\Sigma = r^2 + a^2 \cos^2 \theta, \Delta = r^2 - 2Mr + a^2, A(r) = 1 + r^{-2} \alpha_{22} M^2.$$

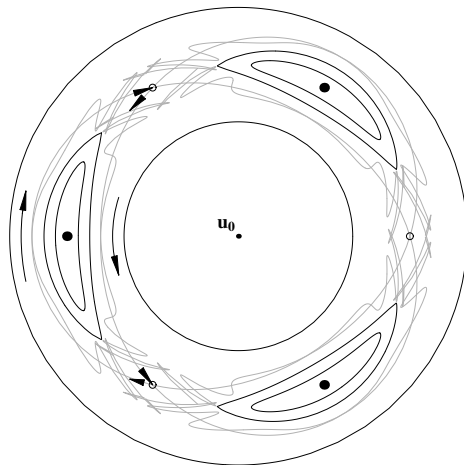
- α_Q controls **integrability**, $\alpha_Q \neq 0 \rightarrow$ **no Carter constant** (“**non-Kerr**”)
- α_{22} affects **frame-dragging**, $\alpha_Q = 0$ and $\alpha_{22} \neq 0 \rightarrow$ **Carter constant** (“**Kerr-like**”)
- $\alpha_Q = \alpha_{22} = 0 \rightarrow$ **Kerr metric**
- **Perturbed quadrupole moment:** $M_2 = -Ma^2 - \frac{1}{3}M^3 \alpha_Q$

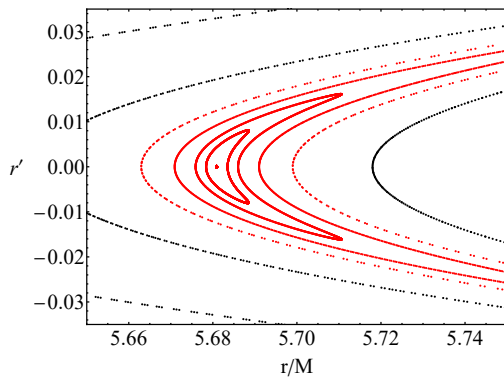
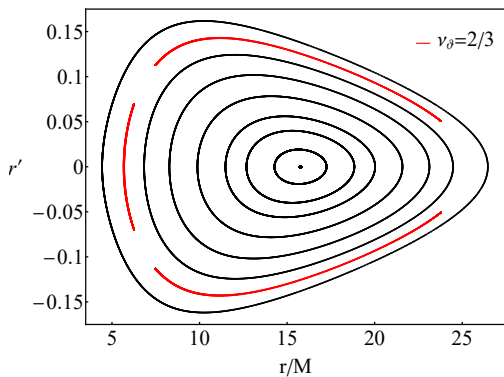
Poincaré map: Signatures of non-integrability

integrability



non-integrability

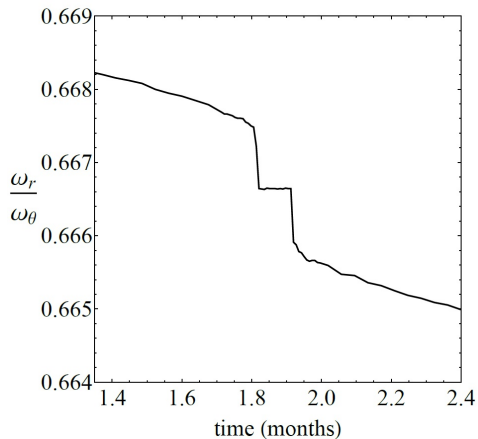




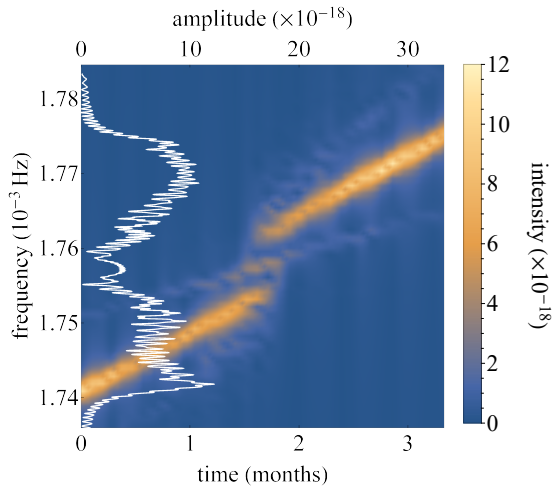
[KD, Kokkotas, GRG (2023)]

- Away from resonances, **the curves surround the central point of the map.**
- **Stable periodic points** are encapsulated by **nested islands.**
- Geodesics share the **same rational ratio $\nu_\theta = \omega_r/\omega_\theta$ throughout the island.**

Signatures of non-integrability in EMRIs



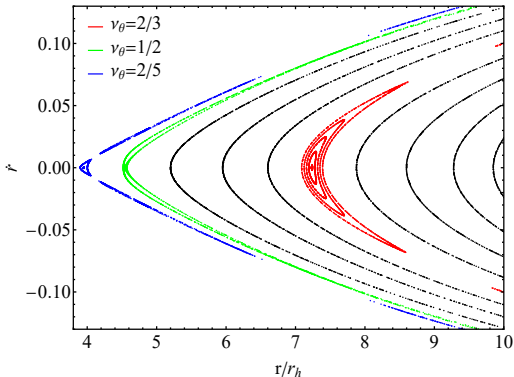
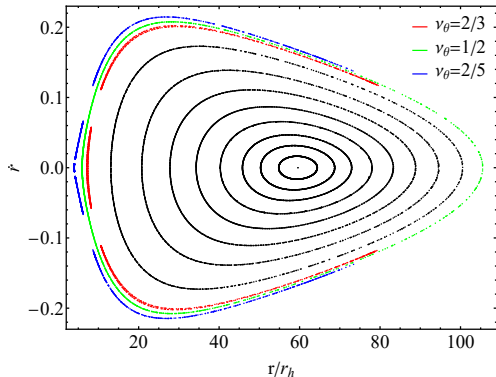
Average crossing time: ~ 300 cycles



[KD, Suvorov, Kokkotas, PRL (2021)]

Outro: Rotating BHs in astrophysical environments and chaos

- **First solution** of a general-relativistic **rotating BH within a matter halo**. [Fernandes, Cardoso, 2507.04389]
- **First geodesic analysis:** General-relativistic, rapidly-rotating BH within a compact matter halo → **environmentally-induced chaos!** [KD, Fernandes, 2508.20191]



[KD, Fernandes, 2508.20191]

Conclusions

- **Thousands of EMRIs are expected** to be detected by LISA. [Gair+, J. Phys. (2017)]
- **Precise modeling of EMRI waveforms is still under construction** due to the perplexity of gravitational self-force. [Barack, CQG (2009)]
- **EMRI modeling approaches** (relativistic, or post-Newtonian) are in full effect and **constantly updated** to include various effects, such as **resonances**. [Katz+, PRD (2021), Ruangsri+, PRD (2014), Berry+, PRD (2016), Sperti+, PRD (2021), Gupta+, PRD (2022)]
- Taking into account **the most dominant resonance** ($2 : 3$) is not enough, **sub-dominant resonances introduce further dephasing**. [Levati, Cardenas-Avendanõ, KD, Pani, PRD (2025)]
- **Non-integrability** is imprinted in **EMRI GWs**. [KD, Suvorov, Kokkotas, PRL (2021), KD, Kokkotas, PRD (2021), KD, Kokkotas, GRG (2023), KD, Huez+, GRG (2023), KD, Pani+, PRD (2023)]
- **Can (non-)integrability de-circularize initially-spherical EMRIs?** [Eleni, KD, Apostolatos, Kokkotas, PRD (2024)] (see talk by A. Eleni on Thursday to find out!)
- **High-precision EMRI GW modeling** requires the inclusion of **resonances** and potential **chaotic phenomena** in order to perform **accurate parameter inference**.

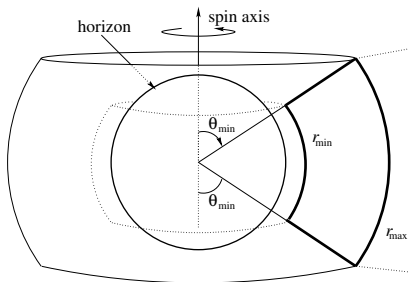


Thanks for your attention!

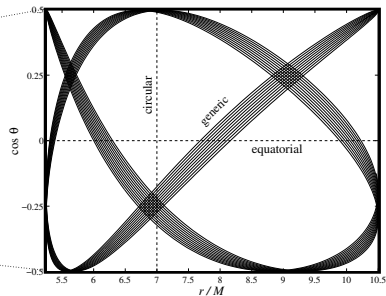
Appendix

LISA targets

- **Extreme-mass-ratio inspirals**
- Supermassive black hole binaries
- Ultra compact binaries in the Milky Way
- Stellar mass black hole binaries years before the merger
- Stochastic gravitational-wave background
- Unforeseen and unmodeled sources



[Drasco, Hughes, PRD (2006)]



Motion is characterized by three orbital frequencies: $\omega_r, \omega_\theta, \omega_\phi$

Initial parameters: $\{\mathcal{E}, \mathcal{L}_z, \vec{x}(0), \dot{\vec{x}}(0)\}$ $\longleftrightarrow$ **Alternative parameters:** $\{e, p, \iota\}$

$$e = \frac{r_{\max} - r_{\min}}{r_{\max} + r_{\min}}, \quad p = \frac{2r_{\max} r_{\min}}{r_{\max} - r_{\min}}, \quad \iota = \frac{\pi}{2} - \theta_{\min}$$

EMRI evolution scheme

The calculation proceeds via the following steps:

1. Define the EMRI initial parameters M , μ , a and α_Q , α_{22}
2. Initialize the trajectory with E_0 , $L_{z,0}$ and $(r(0), \theta(0), \dot{r}(0), \dot{\theta}(0))$
3. Calculate the orbital elements e , p , ι
4. Approximate the fluxes

$$\left\langle \frac{dE}{dt} \right\rangle_{2\text{PN}} = \frac{\mu}{M^2} f_E(e, p, \iota, M_2), \quad \left\langle \frac{dL_z}{dt} \right\rangle_{2\text{PN}} = \frac{\mu}{M} f_L(e, p, \iota, M_2) \quad [\text{Gair, Glampedakis, PRD (2006)}]$$

5. Evolve the coupled (r, θ) geodesic equations with E , L_z changing adiabatically as:

$$E(t) = E_0 + \langle dE/dt \rangle t, \quad L_z(t) = L_{z,0} + \langle dL_z/dt \rangle t$$

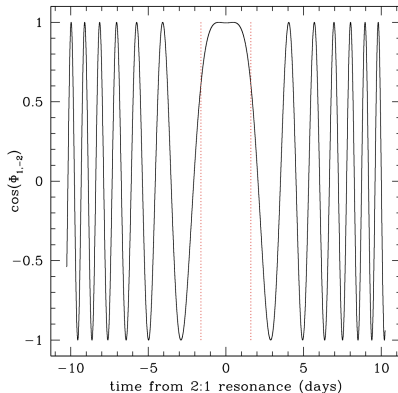
6. Update the fluxes every N_r cycles and repeat from step 2

Phase variables and resonance duration

$$\frac{dJ_k}{d\tau} \sim \dot{J}_k e^{-i\Phi_k(t)}$$

Resonance duration:

$$\begin{aligned}\Phi_k(t) &= \Phi_k(t_{\text{res}}) + (l^*\omega_\theta + m^*\omega_r)(t - t_{\text{res}}) + \frac{1}{2}(l^*\dot{\omega}_\theta + m^*\dot{\omega}_r)(t - t_{\text{res}})^2 + \dots \\ &\simeq \Phi_k(t_{\text{res}}) + \frac{1}{2}(l^*\dot{\omega}_\theta + m^*\dot{\omega}_r)(t - t_{\text{res}})^2.\end{aligned}$$



Resonance	e	ι	p/M	$t_{\text{res}} [\text{d}]$	$T_{\text{max}} [\text{d}]$	$ \mathcal{C}_{\mathcal{E}} $	$ \mathcal{C}_{\mathcal{L}_z} $	$ \mathcal{C}_{\mathcal{Q}} $	$ \mathcal{C}_{\text{tot}} $
Case (i)									
4:3	0.30	0.35	7.46	5.7	362	0.00001	0.00001	0.00006	0.00008
3:2	0.30	0.35	5.36	2.3	84	0.00102	0.00067	0.00208	0.00377
2:1	0.30	0.35	3.59	0.7	10	0.00131	0.00179	0.00046	0.00356
3:1	0.30	0.35	2.92	0.2	1	0.00059	0.00070	0.00310	0.00439
Case (ii)									
4:3	0.30	1.22	11.36	14.6	365	0.00003	0.00004	0.00002	0.00009
3:2	0.30	1.22	8.67	6.5	365	0.00303	0.00123	0.00123	0.00549
2:1	0.30	1.22	6.18	2.0	55	0.00004	0.00080	0.00002	0.00086
3:1	0.30	1.22	5.07	0.5	5	0.00008	0.00024	0.00033	0.00065
Case (iii)									
4:3	0.70	0.35	7.58	11.4	365	0.00001	0.00002	0.00023	0.00026
3:2	0.70	0.35	5.50	4.7	144	0.00127	0.00078	0.00210	0.00415
2:1	0.70	0.35	3.82	1.5	30	0.00167	0.00067	0.00357	0.00591
3:1	0.70	0.35	3.28	0.5	17	0.00026	0.00009	0.00035	0.00070
Case (iv)									
4:3	0.70	1.22	11.47	28.2	365	0.00047	0.00060	0.00002	0.00109
3:2	0.70	1.22	8.80	12.4	365	0.01030	0.00489	0.00261	0.01780
2:1	0.70	1.22	6.36	3.7	88	0.00662	0.00270	0.00494	0.01426
3:1	0.70	1.22	5.41	0.8	5	0.00125	0.00027	0.00126	0.00278

Table 1: Parameter-space location and resonance coefficients for the investigated orbital resonances. We use $a = 0.9$, $\eta = 10^{-5}$ and $M = 10^6 M_{\odot}$. [Flanagan+, PRD (2014), Levati, Cardenas-Avendaño, KD, Pani, PRD (2025)]

Impact of non-integrability: orbital level

w_n : n^{th} crossing of the orbit through a surface of section

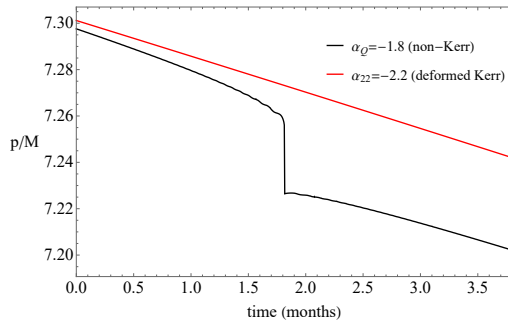
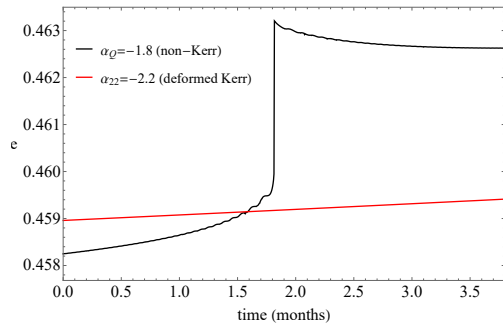
ϑ_n : angle between two successive intersections at w_n, w_{n-1}

Rotation number:

$$\nu_{\theta,N} = \frac{1}{2\pi N} \sum_{i=1}^N \vartheta_i$$

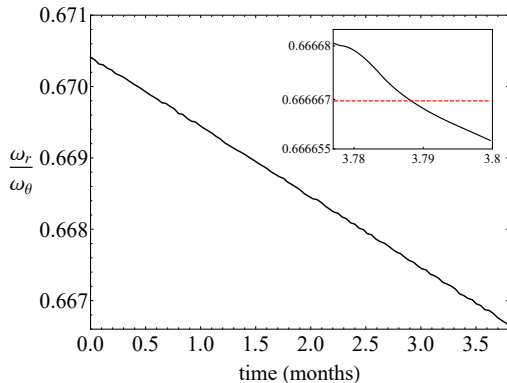
- The limit $N \rightarrow \infty$ converges to $\nu_{\theta} = \omega_r / \omega_{\theta}$ [Contopoulos (2002)]
- Successive rotation numbers form a rotation curve
- The rotation curve of integrable systems is monotonous
- The rotation curve of non-integrable systems exhibits a plateau inside resonant islands

Orbital elements

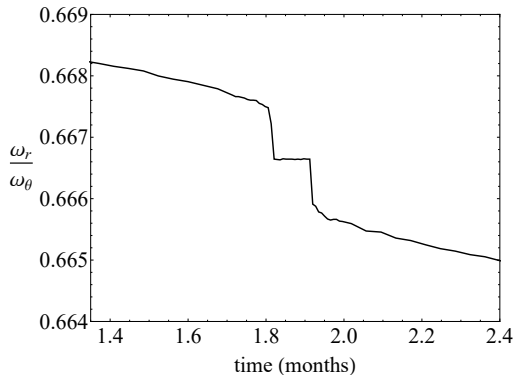


Orbital frequency evolution ($\mu/M = 10^{-6}$)

deformed Kerr

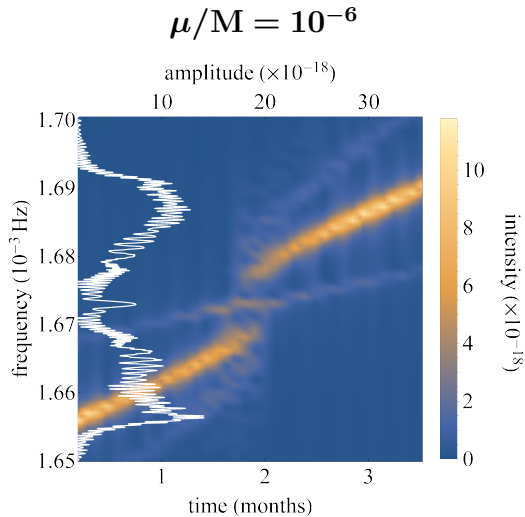
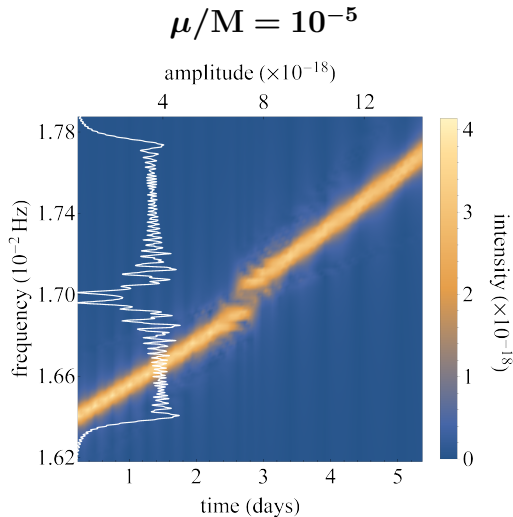


non-Kerr



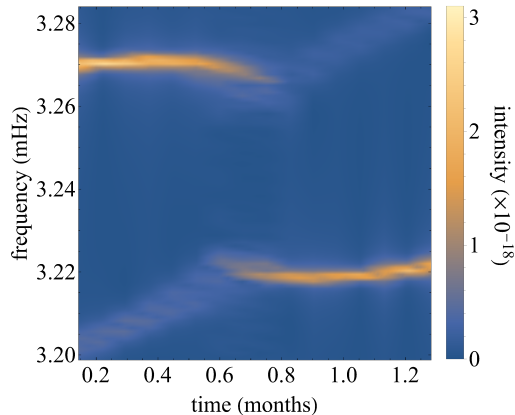
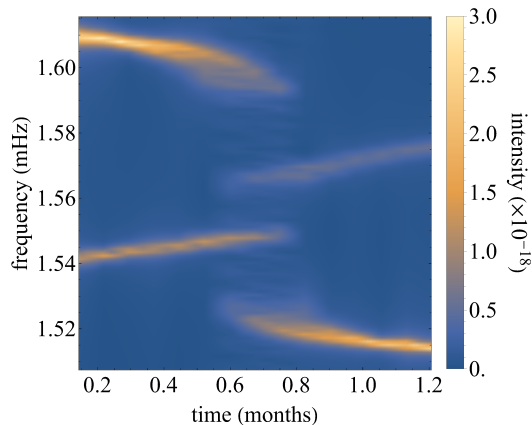
Average time spent in island: ~ 300 cycles

GW glitches in Manko-Novikov EMRIs: a generic feature of chaos



[KD, Kokkotas, PRD (2021)]

EMRIs with supermassive self-interacting rotating boson stars



[KD, Pani+, PRD (2023)]

Waveform modeling and LISA response

Numerical kludge waveforms: Combine exact particle trajectories with approximate expressions for gravitational-wave (GW) emission [Babak et al., PRD (2007)]

- ‘quick and dirty’ EMRI waveforms, perfectly-suited for phenomenology
- remarkable agreement with Teukolsky-based waveforms ($\sim 95\%$)

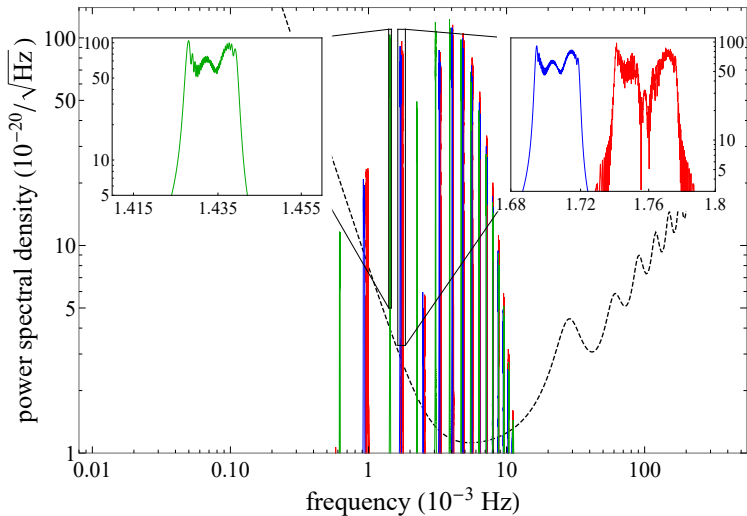
Einstein-quadrupole formula: $h_{ij}^{\text{TT}} = \frac{2}{d} \frac{d^2 I_{ij}(Z^i(t))}{dt^2}$, $\mathbf{Z}(t)$: inspiral trajectory

Incoming GW: $h_{+,\times}(t) = \frac{2\mu}{d} \epsilon_{ij}^{+,\times} \left[\frac{Z^i(t)}{dt^2} Z^j(t) + \frac{Z^i(t)}{dt} \frac{Z^j(t)}{dt} \right]$

LISA response: $h_\alpha(t) \sim [F_\alpha^+(t)h_+(t) + F_\alpha^\times(t)h_\times(t)]$ [Barack, Cutler, PRD (2004)]

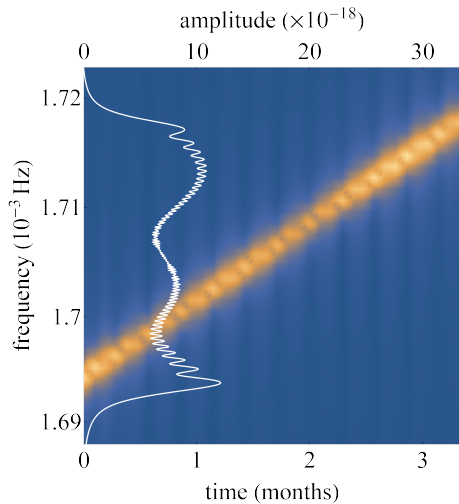
We use a **single-channel approach**, for simplicity, and **omit any noise** in the data stream.

Detectability and GW frequency evolution ($\mu/M = 10^{-6}$)

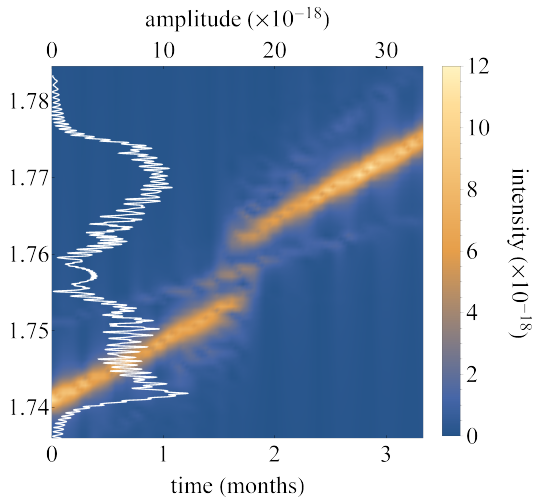


[KD, Suvorov, Kokkotas, PRL (2021)]

deformed Kerr

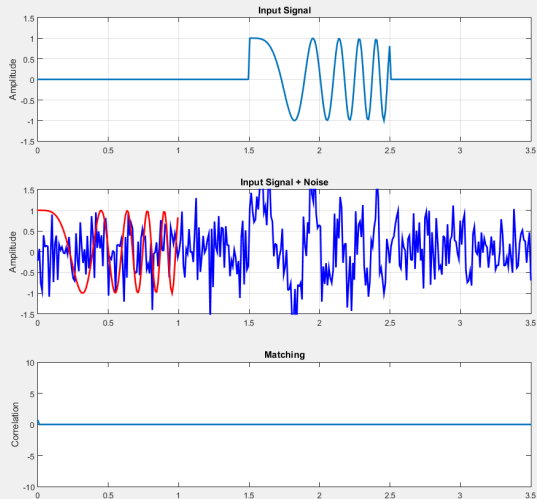


non-Kerr



[KD, Suvorov, Kokkotas, PRL (2021)]

Data analysis: Matched filtering



[Courtesy of: Vitor Cardoso]