

Axisymmetric spacetimes circularity and beyond

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with Jacopo Mazza
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Motivation: Going beyond Kerr

- ✦ In General Relativity the Kerr solution describes vacuum rotating black hole
- ✦ It is important to construct deformations of the Kerr spacetime to have a **benchmark for testing General Relativity** and look for **signatures of modified gravity**.
- ✦ One approach is to find (construct) rotating BH solutions
- ✦ **Alternative approach is to be agnostic and to construct an *ad hoc* metric that describes a deviation from the Kerr metric**

Motivation: Circularity

- ✦ Circularity is accidental symmetry of Kerr solution. A symmetry that is in addition to stationarity and axisymmetry [Also the talk by Cynthia Belen Arias Pruna]
- ✦ In GR, for vacuum solutions circularity is automatic [Even asymptotically "swirling" solution, talk by José Barrientos]
- ✦ **Beyond GR : do we keep circularity or not?** Usually yes, but this is not always justified.

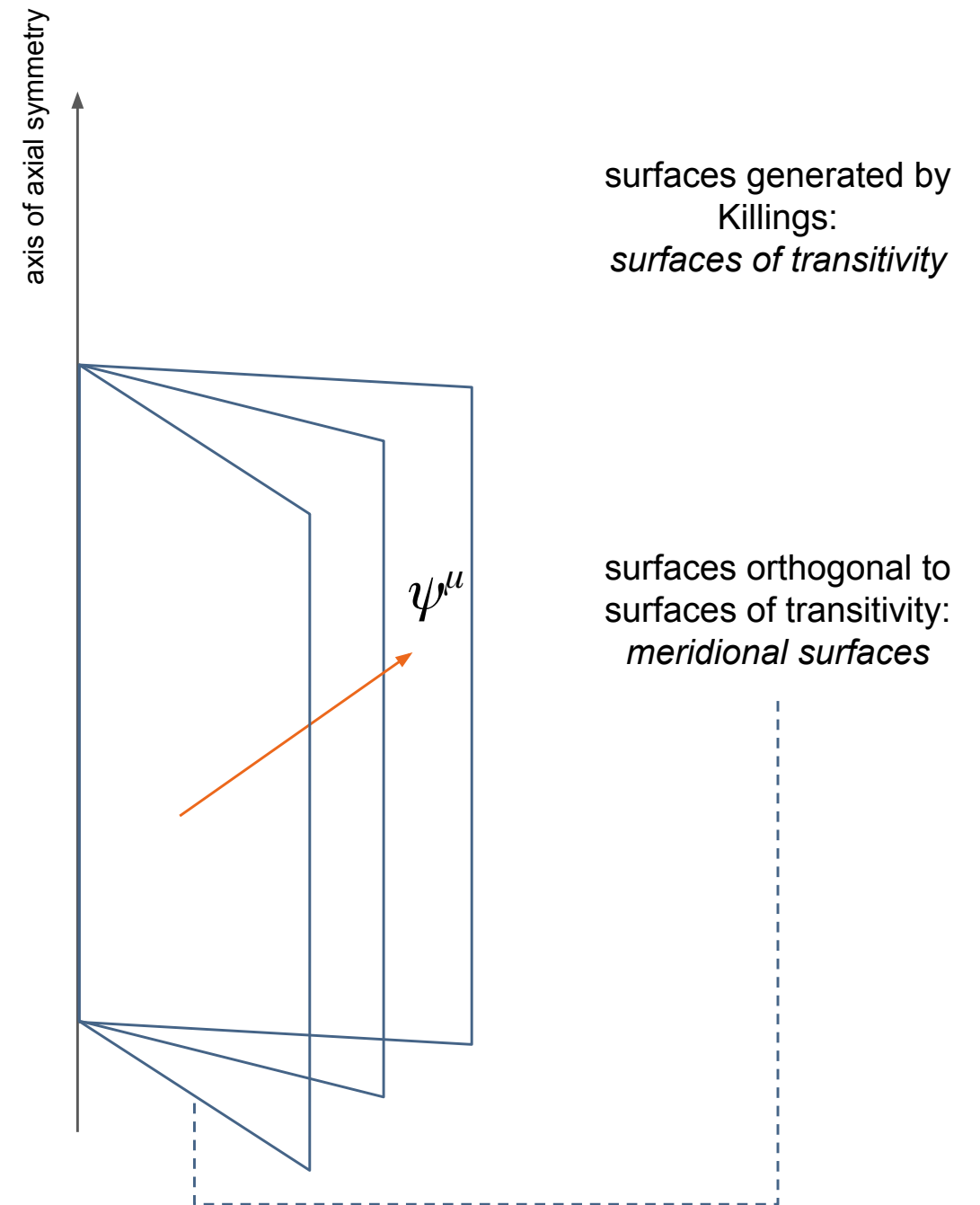
Circularity vs. non-circularity

- ✦ Most general stationary axisymmetric spacetime: two Killing vectors associated to the symmetries
 - ξ^μ (stationarity)
 - ψ^μ (axial symmetry)
- ✦ A spacetime is said to be **circular** if it can be foliated by codimension-2 surfaces orthogonal to the Killing vectors.

by Frobenius' theorem:

$$\xi_{[\mu} \psi_{\nu} \partial_{\rho} \xi_{\sigma]} = 0$$
$$\xi_{[\mu} \psi_{\nu} \partial_{\rho} \psi_{\sigma]} = 0$$

N.B. another definition of circularity involves matter. They are "equivalent".



Carter, Black holes
equilibrium states,
Les Houches lecture notes '73

Circularity vs. non-circularity

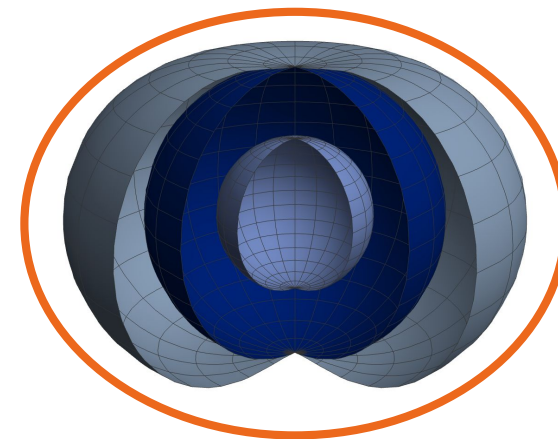
If such foliation exist, one can choose coordinates adapted to foliation

The metric is block diagonal

$$g_{\mu\nu} = \begin{pmatrix} * & 0 & 0 & * \\ 0 & * & * & 0 \\ 0 & * & * & 0 \\ * & 0 & 0 & * \end{pmatrix}$$

$t - \phi$ symmetry
4 functions of two variables

For Kerr metric
Boyer-Lindquist coordinates



Why circularity?

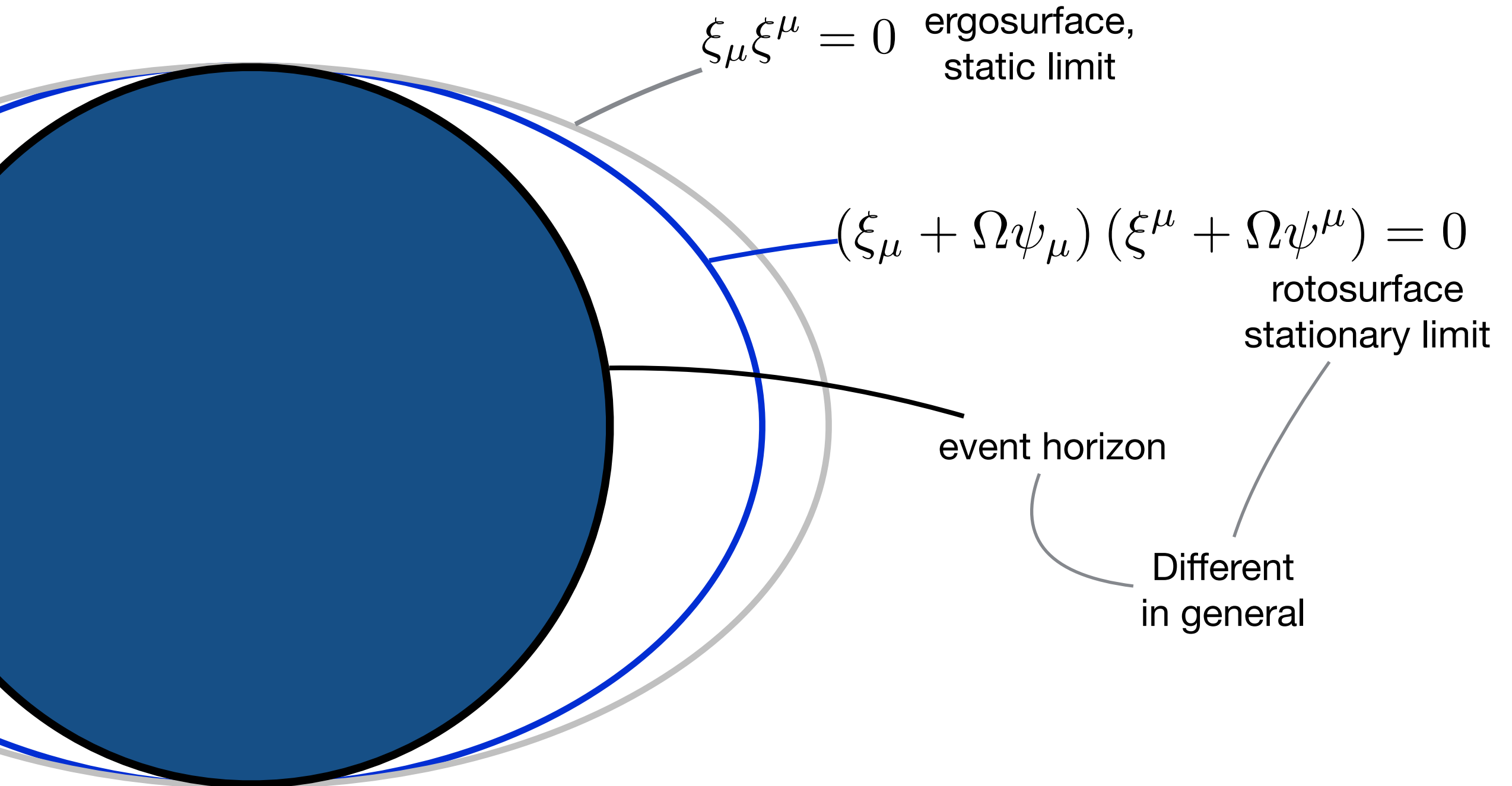
- ⚡ Vacuum GR solutions are circular
- ⚡ Some solutions to alternative theories of gravity are circular:
 - stealth solutions, whereby the metric is that of Kerr;
 - numerical solutions in dilatonic EGB theory;
 - perturbative rotating solutions in sEGB and dynamical Chern-Simons gravity;
- ⚡ beyond-GR solutions that are continuously connected to GR are circular
[Y. Xie et al'21]



Most importantly — Simplicity

- ❖ Metric is block diagonalizable
- ❖ And other features...

Notable surfaces



Circularity means simplicity

If circularity, then:

- ⚡ rotosurface \equiv horizon
- ⚡ frame dragging $\Omega = \text{const.}$ on the horizon (Rigidity theorem)
- ⚡ event horizon is Killing
- ⚡ different notions of surface gravity agree
- ⚡ surface gravity is constant (0th law of BH mechanics)

In addition:

Circularity is necessary (although not sufficient) condition for separability

Circular metrics are useful for phenomenological studies

Why **NON**circularity?

- ✦ Even in GR non-vacuum solutions may have a non-circular metric Gourgoulhon, Bonazzola '93
- ✦ Examples of non-circular solutions:
 - Disformed Kerr black hole Anson, EB, Charmousis, Hassaine '21;
 - Numerical rotating BHs in Einstein-Aether Adam, Figueras, Jacobson, Wiseman '21
- ✦ Circularity is an additional assumption. In modified gravity this assumption translates to the symmetry assumption on extra fields.
 - E.g. in scalar-tensor theory one needs $\xi^\mu \partial_\mu \varphi = \psi^\mu \partial_\mu \varphi = 0$ to fall into class of theories with circular solutions.

Assuming circularity *a priori* is often not justified

What happens when circularity is broken?

Find circularity condition?

The circularity condition are PDEs

$$\xi_{[\mu} \psi_{\nu} \partial_{\rho} \xi_{\sigma]} = 0$$

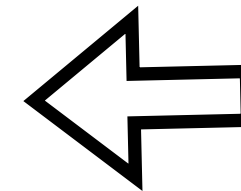
$$\xi_{[\mu} \psi_{\nu} \partial_{\rho} \psi_{\sigma]} = 0$$

Given a metric it is easy
to check circularity

It is difficult to construct
simple non-circular
metrics



Need to "solve" conditions to
get algebraic relations?



Depends on choice
of "good" gauge

Gauge choice for general axisymmetric stationary metric

- ✦ Coordinates adapted to symmetries, i.e. v and ϕ along the Killing orbits,

$$\xi^\mu \partial_\mu = \partial_v \quad \text{and} \quad \psi^\mu \partial_\mu = \partial_\phi$$

Other coordinates r and θ , so that

$$g_{\mu\nu}(x^\alpha) = g_{\mu\nu}(r, \theta).$$

- ✦ These coordinates are defined up to

$$\tilde{v} = v + V(r, \theta)$$

$$\tilde{r} = R(r, \theta)$$

$$\tilde{\theta} = \Theta(r, \theta)$$

$$\tilde{\phi} = \phi + \Phi(r, \theta)$$

$$g_{\mu\nu} = \begin{pmatrix} * & * & * & * \\ & * & * & * \\ & & * & * \\ & & & * \end{pmatrix}$$

where $* = *(r, \theta)$

Most general choice of coords
adapted to the symmetries
(10 components of metric)

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Gauge choice: kill 4 components of the metric?

It is possible locally (using *Cauchy-Kovalevskaya theorem*)

[EB & Mazza'25]

Solving circularity conditions

Kerr-like gauge is great

$$\begin{aligned}\xi_{[\mu}\psi_{\nu}\partial_{\rho}\xi_{\sigma]} &= 0 \\ \xi_{[\mu}\psi_{\nu}\partial_{\rho}\psi_{\sigma]} &= 0\end{aligned}$$

Plug into Kerr
(orthogonal) gauge

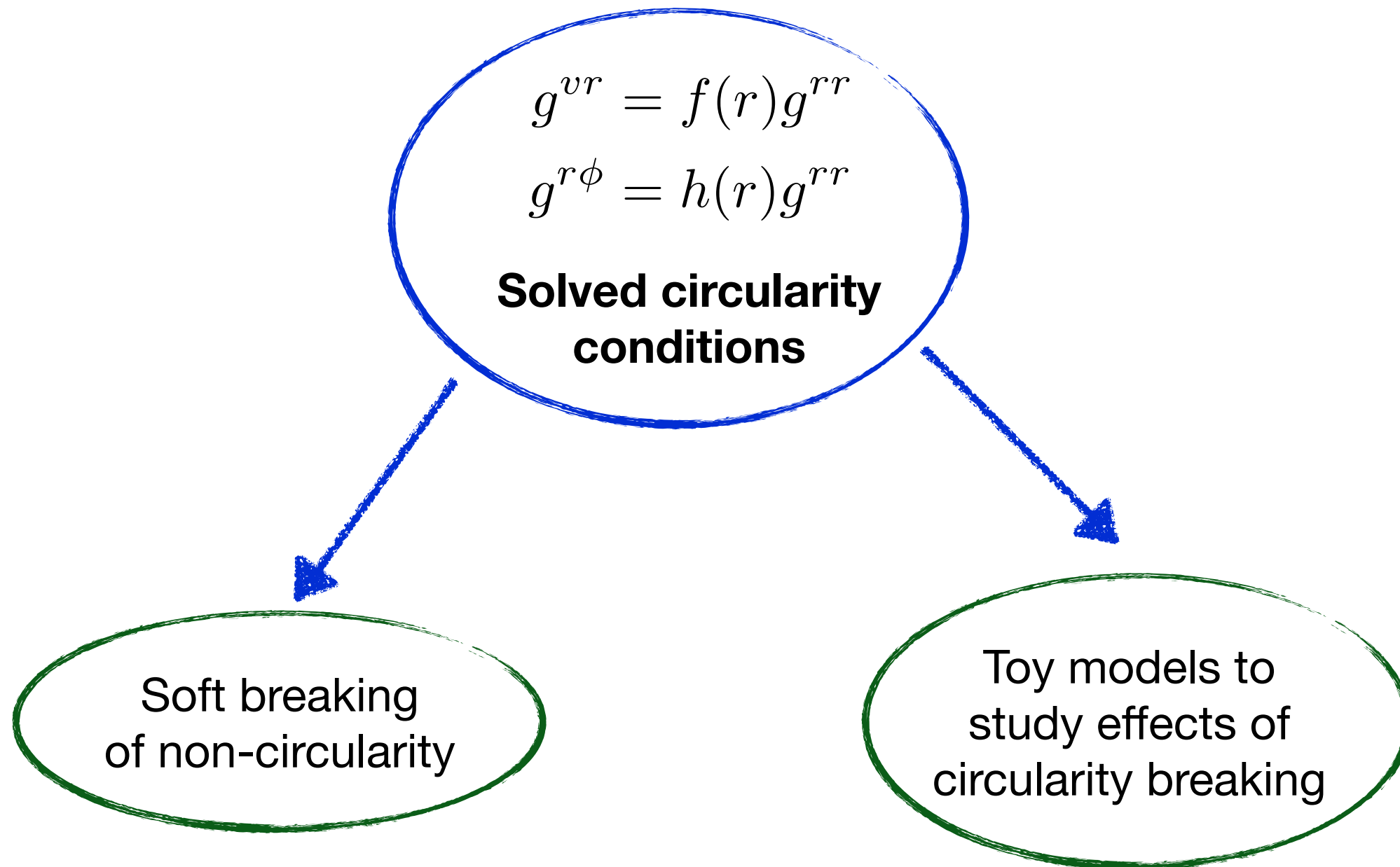
Long calculations



$$\begin{aligned}g^{vr} &= f(r)g^{rr} \\ g^{r\phi} &= h(r)g^{rr}\end{aligned}$$

N.B. The use of the
upper indices !

Why is this interesting?



Examples (see in [arxiv:2505.08880](https://arxiv.org/abs/2505.08880)) of minimal and not-so-minimal deformations of Kerr metric

Conclusions and outlook

- ✦ Gauge choice: orthogonal and Kerr-like gauges can be chosen
 - ✦ Conditions of circularity are solved
 - ✦ Examples
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- ✦ Non-constant surface gravity: Hawking radiation and thermodynamics?
 - ✦ Observational consequences: QNM, geodesics, shadows
 - ✦ Superradiance
 - ✦ Residual gauge: can we simplify further?